

EFFECTS OF HEAT SOURCE, CHEMICAL REACTION, THERMAL DIFFUSION AND MAGNETIC FIELD ON DEMIXING OF A BINARY FLUID MIXTURE FLOWING OVER A STRETCHED SURFACE

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Abstract— Effects of viscous dissipation, heat source, thermal diffusion and chemical reaction on demixing of a binary mixture of incompressible viscous electrically conducting fluids in two dimensional magnetohydrodynamic boundary layer flow on a vertical stretching surface are investigated. The momentum, energy and concentration equations are reduced to non-linear coupled ordinary differential equations by similarity transformations and are solved numerically by using MATLAB's built in solver bvp4c. These numerical results are exhibited graphically from which it has been found that the effects of various parameters are to separate the components of the binary mixture by collecting the lighter and rarer component near the plate and throwing the heavier one away from it.

Keywords— heat and mass transfer, binary fluid mixture, thermal diffusion and chemical reaction, heat source, magnetic field.

I. INTRODUCTION

Separation processes of an electrically conducting binary mixture of incompressible viscous fluids under the influence of magnetic field are considered to be of significant importance due to their applications in many engineering problems such as nuclear reactors and those dealing with liquid metals. A large amount of research work has been done in the field of chemical reaction, heat and mass transfer. There has been renewed interest in studying hydromagnetic flow and heat transfer of continuously stretched surfaces in the presence of a weak transverse magnetic field. This is because hydromagnetic flow and heat transfer have become more important industrially and in many branches of science and engineering.

A study has been carried out to obtain the non-linear MHD flow with heat and mass transfer characteristic of an incompressible viscous, electrically conducting and Boussinesq fluid on a vertical isothermal, cone surface with heat generation/ absorption by El-Kabier *et al.* (2007); Kandasamy and Devi (2004) studied effects of chemical reaction, heat and mass transfer on non-linear laminar boundary layer flow over a wedge with suction and injection. Takhar *et al.* (2000) investigated the flow and mass diffusion of chemical species with first order and higher order re-actions over a continuous stretching sheet with an applied magnetic field.

A study on non-linear hydromagnetic flow, heat and mass transfer over an accelerating vertical surface with internal heat generation and diffusion effects is carried

by Kandasamy and Periasamy (2005). Singh (2001) analyzed the MHD free convection and mass transfer flow with heat source and thermal diffusion. The paper deals with the study of free convection and mass transfer flow of an incompressible, viscous and electrically conducting fluid past a continuously moving infinite vertical plate in the presence of large suction and under the influence of uniform magnetic field considering heat source and thermal diffusion. A study on MHD free convective flow of an incompressible viscous dissipative fluid in an infinite vertical oscillating plate with constant heat flux has been carried by Kishan *et al.* (2006).

The problem of stretching surface with constant surface temperature was analyzed by Afzal (1993). The process involving the mass transfer effect has long been recognized as important principality in chemical processing equipment. Recently, the non linear MHD flow with heat and mass transfer characteristic of an incompressible, viscous, electrically conducting and Boussinesq fluid on a vertical stretching surface with chemical reaction and thermal stratification effects are studied by Kandasamy *et al.* (2005) and Kishan and Amrutha (2010).

Sharma and Singh (2008, 2009, 2010), Sharma and Nath (2012) and Sharma *et al.* (2011, 2012) studied the effect of magnetic field on demixing of a binary fluid mixture. Sharma and Singh (2004, 2007) studied the effect of temperature gradient on demixing of species in hydromagnetic flow of a binary mixture of incompressible viscous fluids between two parallel plates, first taking the plates horizontal and second by taking the plates vertical. They found that the effect of temperature gradient is to separate the components of the binary mixture and the magnetic field increases the effect of species demixing.

The present paper deals with two dimensional steady nonlinear MHD boundary layer flow of an incompressible viscous electrically conducting binary mixture of fluids flowing over a vertical stretching surface in presence of uniform magnetic field by taking into account the viscous dissipation, heat source, chemical reaction and thermal diffusion with a motive to study the effects of various parameters on demixing of the binary fluid mixture.

II FORMULATION OF THE PROBLEM

Consider a two-dimensional steady nonlinear MHD boundary layer flow, heat and mass transfer of a binary fluid mixture of electrically conducting fluids flowing

over a vertical stretching surface in presence of uniform weak magnetic field by taking into account the viscous dissipation with chemical reaction and thermal diffusion effects. The x -axis is taken parallel to the vertical surface and the y -axis is taken normal to it. A transverse weak magnetic field of strength B_0 is applied parallel to y -axis. The fluid properties are assumed to be constant in a limited temperature range. The chemical reactions are taking place in the flow and the physical properties μ , ρ , D and the rate of chemical reaction k_1 are constant throughout the fluid. It is assumed that the induced magnetic field, the external electric field and the electric field due to the polarization of charges are negligible. Under these assumptions, the governing boundary layer equations of momentum, energy and diffusion under Boussinesq approximation are

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \quad (1)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = U \frac{dU}{dx} + v \frac{\partial^2 u}{\partial y^2} + g\beta(T - T_\infty) + g\beta^*(C - C_\infty) - \left(\frac{\sigma B_0^2}{\rho} \right) u, \quad (2)$$

$$\rho c_p \left(u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) = \alpha \frac{\partial^2 T}{\partial y^2} + Q(T - T_\infty) + \mu \left(\frac{\partial u}{\partial y} \right)^2, \quad (3)$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D \left[\frac{\partial^2 C}{\partial y^2} + S_T \frac{\partial C}{\partial y} \frac{\partial T}{\partial y} + S_T C \frac{\partial^2 T}{\partial y^2} \right] - k_1(C - C_\infty), \quad (4)$$

where u and v denote the velocity components in x , y -directions respectively, U is the potential flow velocity, v is the kinematic viscosity, g is the acceleration due to gravity, β is the coefficient of thermal expansion, β^* is the coefficient of concentration expansion, T is the temperature of the fluid mixture, T_∞ is the free stream temperature, T_w is the surface temperature of the plate, C is the concentration of the species, C_w is the concentration of the species at the plate, C_∞ is the concentration of the species far away from the wall, σ is the electrical conductivity of the medium, ρ is the density of the fluid mixture, c_p is the specific heat at constant pressure, α is the thermal conductivity of the fluid mixture, Q is the heat generation coefficient, μ is the coefficient of viscosity, D is the molecular diffusion coefficient, S_T is the Soret number and k_1 is the chemical reaction coefficient.

The boundary conditions are

$$\left. \begin{array}{l} u=U(x)=ax, v=0, T=T_w(x), C=C_w(x) \text{ at } y=0 \\ u \rightarrow 0, T \rightarrow T_\infty(x) = (1-n) T_0 + n T_w(x), \\ C \rightarrow C_\infty \text{ as } y \rightarrow \infty \end{array} \right\} \quad (5)$$

where a is dimensional constant and n is a constant which is the thermal diffusion parameter and is such that $0 \leq n \leq 1$, n being the thermal diffusion parameter is equal to $m_1/(1+m_1)$ of Nakayama and Koyama (1989), where m_1 is constant. T_0 is constant reference temperature say, $T_\infty(0)$. The suffix w and ∞ denote surface and ambient conditions respectively.

Now we introduce similarity variables as follows:

$$\psi = (vxU(x))^{1/2} f(\eta), \quad (6)$$

$$\eta = \left(\frac{U(x)}{vx} \right)^{1/2} y, \quad (7)$$

The velocity components are given by:

$$u = \frac{\partial \psi}{\partial y}, \quad v = -\frac{\partial \psi}{\partial x}. \quad (8)$$

It can be easily verified that the continuity Eq. (1) is identically satisfied. Similarity solutions exist if we assume that $U(x) = ax$ and introduce the non dimensional form of temperature and concentration as:

$$\theta = \frac{(T - T_\infty)}{(T_w - T_\infty)}, \quad \phi = \frac{(C - C_\infty)}{(C_w - C_\infty)}. \quad (9)$$

The set of Eqs. (2)-(4) are coupled partial non-linear differential equations. Introducing the relation (6)-(8) into the Eqs- (2)-(4) we obtain the following non-linear coupled ordinary differential equations:

$$f''' + ff'' + G_r Re \theta + G_c Re \phi - M^2 f' - f'^2 + 1 = 0, \quad (9)$$

$$\theta' + Prf \theta' + Spf \theta + Br f''^2 = 0, \quad (10)$$

$$\phi' + t_d (\phi' \theta' + \phi' \theta'') + S_c f \phi' - S_c \gamma Re \phi = 0, \quad (11)$$

where $G_r = vg\beta^* (T_w - T_\infty)/U^3$ is the Grashof number, $Re = Ux/v$ is the Reynolds number, $G_c = vg\beta^* (C_w - C_\infty)/U^3$ is the modified Grashof number, $M^2 = \sigma B_0^2 / (\rho a)$ is the magnetic parameter, $Pr = \mu c_p / \alpha$ is the Prandtl number, $S = Q / (\rho a c_p)$ is the heat source parameter, $Br = \mu U^2 / (\alpha (T_w - T_\infty))$ is the Brinkman number, $S_c = v/D$ is the Schimdt number, $\gamma = v k_1 / U^2$ is the chemical reaction parameter and $t_d = S_T (T_w - T_\infty)$ is the thermal diffusion number.

The boundary conditions given by Eqs. (5), in light of Eqs. (6) to (8), become

$$\left. \begin{array}{ll} f=0, f'=1, \theta=1, \phi=1 & \text{at } \eta=0 \\ f' \rightarrow 0, \theta \rightarrow 0, \phi \rightarrow 0 & \text{as } \eta \rightarrow \infty \end{array} \right\} \quad (12)$$

Since the solutions of the set of non-linear coupled ordinary differential equations (9)-(11) under the boundary conditions (12) cannot be obtained in a closed form therefore we have solved these equations numerically with MATLAB's built-in solver bvp4c.

III RESULTS AND DISCUSSION

In order to get a clear insight of the physical problem, numerical calculations have been carried out for various values of the parameters M , G_r , G_c , Re , S , γ , S_c and t_d , and these numerical results for concentration of the lighter component of the binary fluid mixture are plotted against η for various values of above mentioned parameters and are displayed in Figs. 1-8. It is observed from the figures that the concentration of the lighter component of the binary mixture is more at the plate and decreases exponentially as η increases from 0 to 2.

Figure 1 and 8 depict that the concentration of the lighter and rarer component of the binary mixture increases with increase in the values of the parameters M and t_d and from Figs. 2-7 reverse effect is observed with increase in the values of the parameters G_r , G_c , Re , S , γ and S_c .

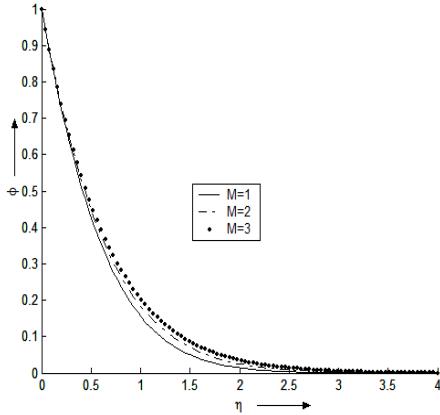


Figure 1. The graph of ϕ against η , $G_r=1$, $Re=3$, $G_c=3$, $Pr=0.71$, $S=0.2$, $Br=2$, $S_c=0.62$, $\gamma=1$ and $t_d=0.001$ for various values of M .

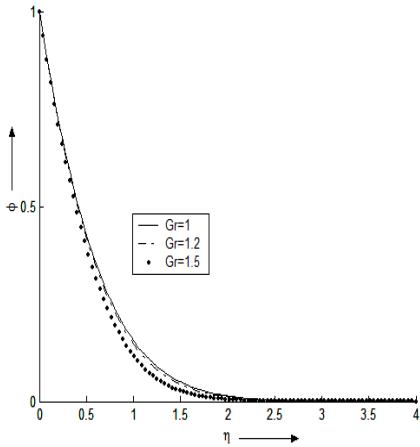


Figure 2. The graph of ϕ against η , $M=1$, $Re=3$, $G_c=3$, $Pr=0.71$, $S=0.2$, $Br=2$, $S_c=0.62$, $\gamma=1$ and $t_d=0.001$ for various values of G_r .

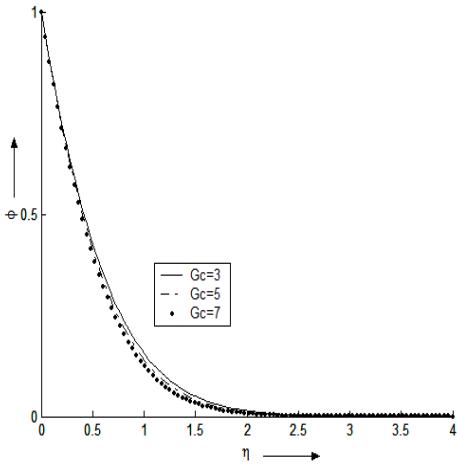


Figure 3. The graph of ϕ against η , $G_r=1$, $Re=3$, $M=1$, $Pr=0.71$, $S=0.2$, $Br=2$, $S_c=0.62$, $\gamma=1$ and $t_d=0.001$ for various values of G_c .

Thus we conclude that the effects of all these parameters is to demix the binary mixture by collecting the rarer and lighter component of the binary fluid mixture near the plate and throwing the heavier component

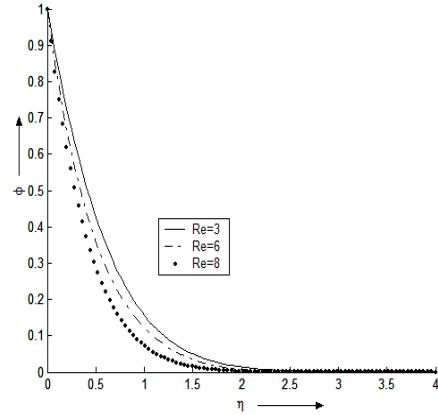


Figure 4. The graph of ϕ against η , $M=1$, $G_r=1$, $G_c=3$, $Pr=0.71$, $S=0.2$, $Br=2$, $S_c=0.62$, $\gamma=1$ and $t_d=0.001$ for various values of Re .

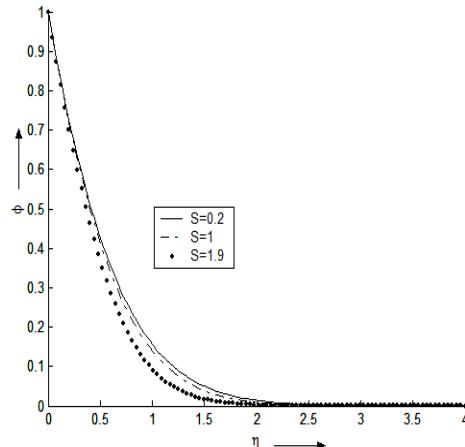


Figure 5. The graph of ϕ against η , $G_r=1$, $Re=3$, $G_c=3$, $Pr=0.71$, $M=1$, $Br=2$, $S_c=0.62$, $\gamma=1$ and $t_d=0.001$ for various values of S .

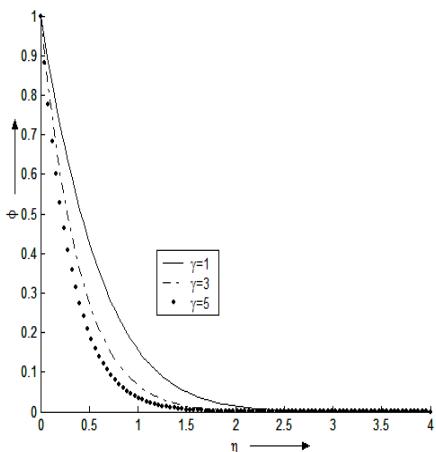
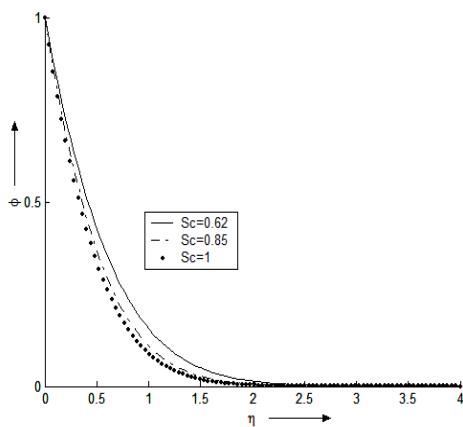
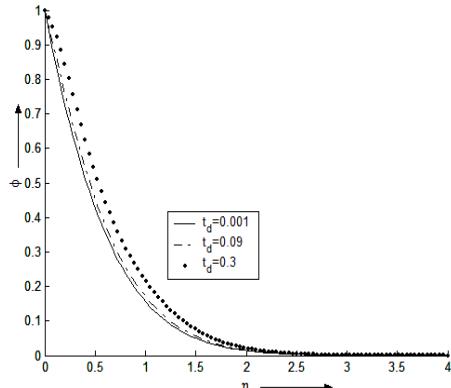


Figure 6. The graph of ϕ against η , $M=1$, $G_r=1$, $Re=3$, $G_c=3$, $Pr=0.71$, $S=0.2$, $Br=2$, $S_c=0.62$ and $t_d=0.001$ for various values of γ .

away from it. From the process of numerical computation, the local skin friction, the Nusselt number and the Sherwood number, which are respectively proportional to $f''(0)$, $\theta'(0)$ and $\phi'(0)$. are also worked out and their numerical values are presented in a tabular form.

Table 1: Numerical values of $f''(0)$, $\theta'(0)$ and $\phi'(0)$ for $Pr=0.71$ and $Br=2$.

M	G_r	G_c	Re	S	γ	t_d	S_c	$f''(0)$	$\theta'(0)$	$\phi'(0)$
1	1	3	3	0.2	1	0.001	0.62	3.9082	4.1885	-1.5392
2	1	3	3	0.2	1	0.001	0.62	1.7357	0.4978	-1.4958
3	1	3	3	0.2	1	0.001	0.62	-0.0073	0.0699	-1.4648
1	1	3	3	0.2	1	0.001	0.62	3.9082	4.1885	-1.5392
1	1.2	3	3	0.2	1	0.001	0.62	4.6562	6.4519	-1.5530
1	1.5	3	3	0.2	1	0.001	0.62	7.3881	16.1273	-1.6021
1	1	3	3	0.2	1	0.001	0.62	3.9082	4.1885	-1.5392
1	1	5	3	0.2	1	0.001	0.62	6.0924	10.1691	-1.5659
1	1	7	3	0.2	1	0.001	0.62	8.3584	18.8190	-1.5888
1	1	3	3	0.2	1	0.001	0.62	3.9082	4.1885	-1.5392
1	1	3	6	0.2	1	0.001	0.62	0.8616	-16.8620	-2.0059
1	1	3	8	0.2	1	0.001	0.62	5.9991	-1.0752	-2.3286
1	1	3	3	0.2	1	0.001	0.62	3.9082	4.1885	-1.5392
1	1	3	3	1	1	0.001	0.62	5.1997	10.6317	-1.5658
1	1	3	3	1.9	1	0.001	0.62	11.0830	55.7155	-1.6346
1	1	3	3	0.2	1	0.001	0.62	3.9082	4.1885	-1.5392
1	1	3	3	0.2	3	0.001	0.62	3.1397	2.1950	-2.4478
1	1	3	3	0.2	5	0.001	0.62	2.7824	1.4925	-3.1109
1	1	3	3	0.2	1	0.001	0.62	3.9082	4.1885	-1.5392
1	1	3	3	0.2	1	0.09	0.62	4.0738	4.6545	-1.1559
1	1	3	3	0.2	1	0.3	0.62	4.5509	6.1686	-0.4557
1	1	3	3	0.2	1	0.001	0.62	3.9082	4.1885	-1.5392
1	1	3	3	0.2	1	0.001	0.85	3.5893	3.2686	-1.7922
1	1	3	3	0.2	1	0.001	1	3.4342	2.8621	-1.9387

Figure 7. The graph of ϕ against η , $M=1$, $G_r=1$, $Re=3$, $G_c=3$, $Pr=0.71$, $S=0.2$, $Br=2$, $\gamma=1$ and $t_d=0.001$ for various values of S_c .Figure 8. The graph of ϕ against η , $M=1$, $G_r=1$, $Re=3$, $G_c=3$, $Pr=0.71$, $S=0.2$, $Br=2$, $S_c=0.62$ and $\gamma=1$ for various values of t_d .

Finally, the effects of the local skin friction, the Nusselt number and the Sherwood number are shown in Table 1. The behaviour of these parameters is self – evident from the Table 1 and hence any further discussion about them seems to be redundant.

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